

## Light Scattering in Media

- When a photon interacts with a particle in a media, it can either be absorbed or reflected.
- Let  $\sigma_a$  be the absorption coefficient and  $\sigma_s$  be the scattering coefficient, then the change of radiance due to out-scattering (loss) is:

$$(\vec{\omega} \cdot \nabla)L(x, \vec{\omega}) = -(\sigma_s + \sigma_a)L(x, \vec{\omega})$$

3



## Light Scattering – Cont'd

- There is also a gain (in-scattering), which is integrated from all directions.
- The gain from in-scattering:

$$(\vec{\omega} \cdot \nabla)L(x, \vec{\omega}) = \sigma_s(x) \int_{\Omega} p(x, \vec{\omega}', \vec{\omega}) L_i(x, \vec{\omega}') d\vec{\omega}'$$

$p$  is the distribution function of the scattered light, called **phase function**.

- The gain from emission (in cases like flames):

$$(\vec{\omega} \cdot \nabla)L(x, \vec{\omega}) = \sigma_a L_e(x, \vec{\omega})$$

4



## The Volume Rendering Equation

- By Integrating the equation (somehow..) on both sides of a segment  $s$ , we get the *volume rendering equation*:

$$L(x, \vec{\omega}) = \int_0^s e^{-\tau(x, x')} \sigma_a(x') L_e(x', \vec{\omega}) dx' +$$

$$\int_0^s e^{-\tau(x, x')} \sigma_s(x') \int_{\Omega} p(x', \vec{\omega}', \vec{\omega}) L_i(x', \vec{\omega}') d\vec{\omega}' dx' +$$

$$e^{-\tau(x, x+s\vec{\omega})} L(x+s\vec{\omega}, \vec{\omega})$$

- $\tau(x, x') = \int_x^{x'} \sigma_t(x'') dx''$  is called the *optical depth*.

5



## The Phase Function

- Describes the distribution of the scattered light.
- Unlike BRDF, it is normalized in the sphere.
  - The actual amount of absorption is controlled by an external coefficient.

$$\int_{\Omega} p(x, \vec{\omega}, \vec{\omega}') d\vec{\omega}' = 1$$

- Often, the function only depends on the angle  $\theta$  between the incoming ray and the scattered ray, and we use  $p(\theta)$ .
  - $\theta=0$  means forward scattering, and  $\theta=\pi$  means backward scattering.

6



## The Phase Function – Cont'd

- Isotropic Scattering:

$$p(\theta) = \frac{1}{4\pi}$$

- A photon will scatter in a random direction

- The Henyey-Greenstein phase function:

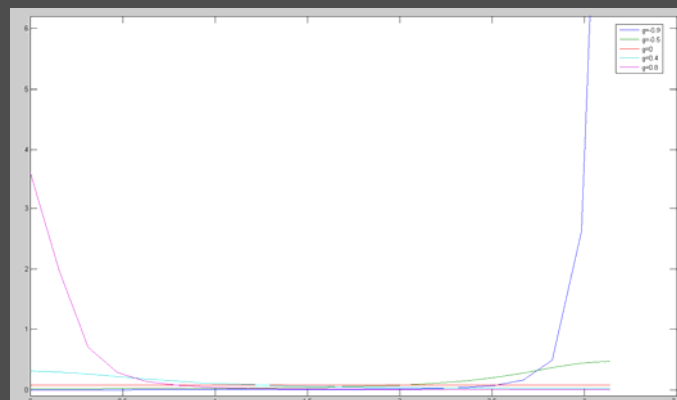
$$p(\theta) = \frac{1 - g^2}{4\pi(1 + g^2 - 2g \cos \theta)^{1.5}}, g \in (-1, 1)$$

7



## HG Phase Function

- $g$  controls the scattering distribution.
  - $g=0$  – isotropic
  - $g>0$  – more forward scattering
  - $g<0$  – more backward scattering



8



## Phase Functions – cont'd

- ❑ The HG phase function has been found to model well scattering in skin, ocean, stone and etc.
- ❑ It is possible to use a linear combination of weighted HG functions for better realistic results.
- ❑ The Schlick phase function approximates HG, without the costly 1.5 power:

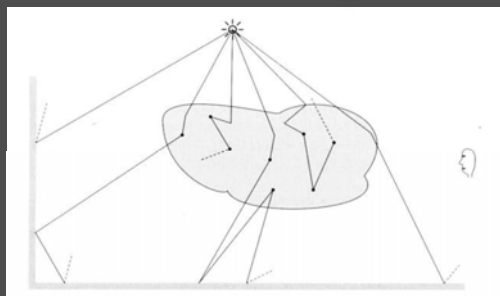
$$p(\theta) = \frac{1 - g^2}{4\pi(1 + k \cos \theta)^2}$$

9



## Photon mapping in Volumes

- ❑ The algorithm is similar, with a few key adaptations.
- ❑ The photon mapping algorithm:
  - Photon tracing
  - Volume radiance estimation
  - Rendering



10



## Creating The Photon Map

- When the photon enters the media (volume), it has an average distance before it forms an intersection:

$$d = \frac{1}{\sigma_t}$$

- Photon scattering is done according to Russian roulette:

$$x \in [0, 1] \rightarrow \begin{cases} x \leq \frac{\sigma_s}{\sigma_t} \Rightarrow \text{scattered} \\ x > \frac{\sigma_s}{\sigma_t} \Rightarrow \text{absorbed} \end{cases}$$

- Color bands are scaled according to the regular algorithm.

11



## Photon Map – Cont'd

- Photon Storing is done in a specified volume map, because radiance estimation is different in volume.
  - A good optimization is to save only photons that scattered at least once.
- Photon can be emitted from the volume with a spectrum based on local conditions.



12



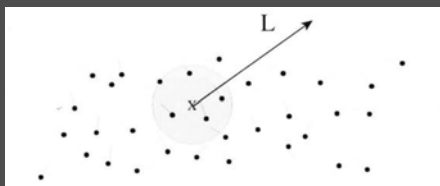
## Radiance Estimate

- The two dimensional approximation does not hold in volumes.

- We estimate the required integral:

$$\begin{aligned}
 (\vec{\omega} \cdot \nabla) L_o(x, \vec{\omega}) &= \sigma_s(x) \int_{\Omega} p(x, \vec{\omega}', \vec{\omega}) L_i(x, \vec{\omega}') d\vec{\omega}' = \\
 &= \int_{\Omega} p(x, \vec{\omega}', \vec{\omega}) \frac{d^2 \Phi(x, \vec{\omega})}{dV} \approx \sum_{i=0}^n p(x, \vec{\omega}_p, \vec{\omega}) \frac{\Phi_p(x, \vec{\omega}_p)}{\frac{4}{3} \pi r^3}
 \end{aligned}$$

In a growing radius  $r$  of a sphere, to engulf  $n$  nearest photons.



13



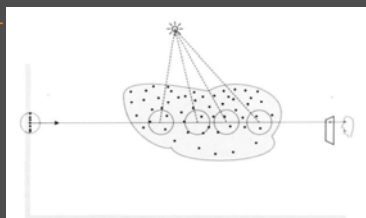
## Rendering – Ray Marching

- Idea: Shoot ray through medium. Divide ray into small segments  $\Delta x$ , and accumulate radiance.
- Radiance is due to single scattering (direct illumination) and multiple scattering (using radiance estimate)
- In the next segment  $n+1$  we get:

$$L_{n+1}(x, \vec{\omega}) = \sum_{i=0}^N L_i(x, \vec{\omega}) p(x, \vec{\omega}_i, \vec{\omega}) \sigma_s(x) \Delta x +$$

$$\sum_{i=0}^n p(x, \vec{\omega}_p, \vec{\omega}) \frac{\Phi_p(x, \vec{\omega}_p)}{\frac{4}{3} \pi r^3} \Delta x +$$

$$e^{-\sigma_t(x) \Delta x} L_n(x + \Delta x, \vec{\omega})$$





- The direct contribution of the light must be attenuated with the distance from the points.

14



Examples



From: <http://graphics.stanford.edu/courses/cs348b-competition/cs348b-04/underwater/index.html>

15 